

## AN EMPIRICAL HBT LARGE SIGNAL MODEL FOR CAD

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**Abstract** — Extensive measurements were performed at different temperatures on GaAs HBT's in order to obtain a simple Large Signal Model usable in CAD. The model was experimentally evaluated by DC, S and Power Spectrum measurements and good correspondence was obtained between measurements and experiment.

## I. INTRODUCTION

Heterojunction Bipolar Transistors (HBT) have become very promising devices for different applications at microwaves and millimeter wave frequencies. An important condition for any successful design work is availability of an accurate large signal model (LSM). In recent years more than 100 publications in respectable journals were devoted on the creation of a compact model for use in CAD tools and procedures for extraction of parameters for these models [1-26]. Despite the tremendous work already done on the subject we still do not have unified, accurate model standard in industry together with the automatic extraction procedure. The reason for this is that the device physics is very complicated, the range of currents is broad, and power density at which the device operates is very high. All this together with poor thermal conductivity of the III-V materials makes the problem of the creation of universal LSM for HBT more difficult. Many of the existing HBT models are based on solid physical background, but again because of the difficulties of the problem they end up with many empirical coefficients, that are difficult to extract.

In some cases the HBT models can be "excessively" accurate and include effects that may not be so important and model simplification is appropriate in such cases. When the model is complicated, an additional difficulty to the extraction problems is that such a model could show problems with the convergence.

## II. DEVICE MODELING

The HBT used in this study are AlGaAs/GaAs HBTs with a 4 x 20  $\mu\text{m}^2$  emitter [27].

Systematic DC, S- multi-bias measurements were performed at different temperatures (-25-+125C) in order to obtain a general picture of the device behavior.

ICCAP (Agilent) was used in the measurement and extraction of basic parameters.

Figures 1-3 shows some results for  $I_b$ ,  $I_{ce}$ , and Forward Gummel (FG)  $\beta$  obtained from these devices that are typical for many HBT. As can be seen, when keeping the  $I_b$  (or  $I_e$ ) constant the voltage shift of the BE junction is almost linear and can be used to monitor the device temperature. Increasing the temperature will decrease the base voltage required to sustain the same current and the temperature coefficient  $T_c V_{bj}$  is ~-0.001. A problem with this type of transistors is that  $\beta$  cannot be considered constant, Fig. 3. The models considering  $\beta$  constant like Gummel-Poon model produce significant error and separate equations for  $I_b$  and  $I_{ce}$  should be used in the model. There is a change of the  $V_{be}$  at which we have maximum of  $\beta$  and this should be considered in the LSM.

The transistor can be described with a conventional equivalent circuit shown in Fig.5. Nonlinear are current sources  $I_{ce}$  and  $C_{be}$  and  $C_{bc}$ , the remaining elements can be considered linear and there is a significant amount of papers describing the extraction of the small signal equivalent circuit [14-26]. For devices exhibiting substrate effects an extra port should be added.

In order to use the fact that when keeping  $I_b$  constant the junction voltages varies linearly with the temperature, we can modify diode current definition. It is a common practice to describe the complicated  $I_b$  dependence by several diodes (respective exponential functions) in order to improve the accuracy and describe the different physical phenomena occurring in the device [1-12]. According to the device physics we use an exponential function to describe the semiconductor junction, but the reference point is changed to  $V_{bj}$  above the knee voltage at currents where we normally operate the device. In order to have a possibility to fit a variety of factors and effects influencing the junction current, we can describe the  $I_b$  as a power series:

$$I_b = I_{be} + I_{bc}$$

$$I_{be} = I_{jbe}(\exp(P_{be}) - \exp(P_{be0})), (1)$$

$$I_{bc} = I_{jbc}(\exp(P_{bc}) - \exp(P_{bc0})), (2)$$

$$P_{be} = P_{be1}(V_{be} - V_{je}) + P_{be2}(V_{be} - V_{je})^2 \dots$$

$$P_{bc} = P_{bc1}(V_{bc} - V_{jc}) + P_{bc2}(V_{bc} - V_{jc})^2 \dots \quad (3)$$

$$P_{be0} = -P_{be1}V_{je} + P_{be2}V_{je}^2 \dots$$

$$P_{bc0} = -P_{bc1}(V_{jc}) + P_{bc2}(-V_{jc}^2) \dots \quad (4)$$

This definition is equivalent to the conventional diode equation, but describes the currents with highest accuracy in the useful operating range, because extraction point is changed. When  $V_{be} = V_{je}$  the base current  $I_{be}$  is equal to  $I_{jbe}$  and at  $V_{be}=0$ ,  $I_{be}=0$ . Usually 1-3 terms of the power series are enough to obtain accuracy 2%. If higher accuracy is required more terms can be added to improve the fit at very small currents.

Using Eq.1-4 we can reduce the number of the diodes (exponents) and keep the accuracy. This in turn will improve the convergence of the large signal model. It is a common practice to limit the exponential growth of the current using a modified exponent (soft exponent). In this case this is not required any more, because we can limit the exponential growth of the current by selecting the coefficients of the arguments of Eq.3 in a proper way.

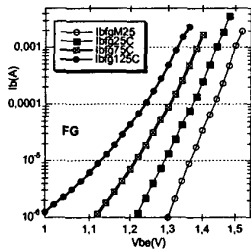


Fig.1. Ib FG

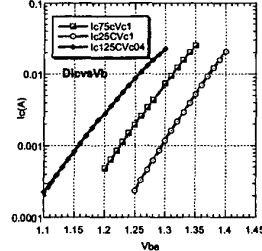


Fig.2. Ic vs. Vbe

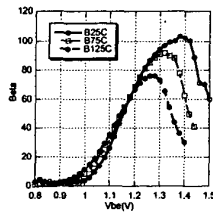


Fig.3. FG β vs. T

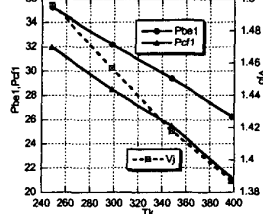


Fig.4. Vjb, Pbe1, Pcf1 vs. T

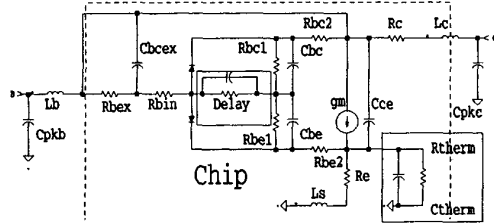


Fig.5. Equivalent circuit

In Fig. 4. and table 1 are shown extracted parameters of the base part using one term in the power series of Eq. 1-3. The changes of the parameters are linear with respect the temperature and this can be used to model self-heating in a simple way using linear dependencies.

We can keep the basic form of the Gummel-Poon model [10] but define collector current source as a power series:

$$I_c = I_{pkc} (\exp(P_{cf}) - \exp(P_{cr})) \left(1 - \frac{V_{be}}{V_{ar}} - \frac{V_{bc}}{V_{af}}\right) \quad (5)$$

$P_{cf}$  is a power series centered at  $V_{bep}$  and with a variable  $V_{be}$ , i.e.:

$$P_{cf} = P_{cf1}(V_{be} - V_{bep}) + P_{cf2}(V_{be} - V_{bep})^2 \dots \quad (6A)$$

$$P_{cr} = P_{cr1}(V_{bc} - V_{bcp}) + P_{cr2}(V_{bc} - V_{bcp})^2 \dots \quad (6B)$$

$$V_{bcp} = V_{bep} + \Delta V_{bep} * (1 + \tanh(P_{cf1} * V_{ce}))$$

When  $V_{be} = V_{bep}$  the collector current is equal to  $I_{pkc}$  and the transconductance is equal to  $g_m = I_{pkc} P_{c1}$ .

Typically with 3 terms of the power series the error is less than 2%. There is freedom for the selection of  $V_{bep}$  and  $I_{pkc}$ , but probably the best is to select  $V_{bep}$  equal to the voltage at which  $\beta$  is maximum and usually even a simple fitting program will converge and find the optimum values.

Normally the accuracy at negative  $V_{ce}$  voltages is not so critical and corresponding parameters from Eq.6A. and 6.B can be considered equal.

The collector current  $I_{ce}$  can be expressed also as:

$$I_c = I_{pkc} (\exp(P_{cf}) \cdot \tanh(\alpha V_{ce}) (1 + \lambda V_{ce})) \quad (5a)$$

$$\alpha = \alpha_r + \alpha_s (\exp(P_{cf1} V_{ce}) - 1)$$

This definition will behave better from numerical point of view, because  $\tanh$  is limited function and at  $V_{ce}=0$ ,  $I_{ce}=0$ .

The exponent's argument  $P_{cf}$  varies with collector voltage for number of reasons - change of the carrier velocity, temperature... This can be reflected using following approach:

$$P_{c1m} = P_{c1} (1 + B_1 / \cosh^2(B_2 V_{ce})) \quad (7)$$

The coefficient  $B_1$  ( $B_1 \approx 0.1-0.3$ ) is extracted at low collector voltages and  $B_2 \approx 10-30$  is a fitting coefficient, which will determine how quickly the gain is changing when  $V_{ce}$  is increased.

### HBT Capacitances.

Devices can operate at high currents and voltages and their nonlinear capacitances should include both depletion and diffusion parts. The total  $C_{be}$  capacitance shows rapid increase and then decrease at high bias. That is similar to what can be found in homo-junction transistors [7,14], but in HBT this increase will be larger than the increase we observe in the homo-junction.

Diffusion capacitances are described as [1]:

$$C_{bedif} = C_{bep} + C_{be0} \cdot (1 + \tanh[C_{be10} + C_{be11} \cdot V_{be}]), \quad (8)$$

$$C_{bcdif} = C_{bcp} + C_{bc0} \cdot (1 + \tanh[C_{bc20} + C_{bc21} \cdot V_{bc}]), \quad (9)$$

Integrating with the terminal voltages we can obtain the  $Q_{be}$  and  $Q_{bc}$  when using the charge definition:

$$Q_{bedif} = C_{be0} \frac{\log[\cosh[C_{be10} + C_{be11} \cdot V_{be}]]}{C_{be11}} + (C_{bep} + C_{be0}) V_{be}$$

$$Q_{bcdif} = C_{bc0} \frac{\log[\cosh[C_{bc20} + C_{bc21} \cdot V_{bc}]]}{C_{bc11}} + (C_{bcp} + C_{bc0}) V_{bc}$$

(10)

The ordinary equation for the depletion part can create convergence problems when  $V_{bc} = V_{bci}$  and to avoid this the equation for the depletion capacitance equation is modified to:

$$C_{dep} = C_{depp} + C_{dep0} (x^2 + m)^{-n-1} (m - (2n-1)x^2) \quad (11)$$

where  $x = V_{bc} - V_{bci}$ ,  $n$ - grading coefficient,  $m$ - parameter determines the maximum- minimum capacitance ratio ( $0 < m < 0.5$ ). Eq.11 gives results identical to the textbooks depletion capacitance definition, but it is well defined at  $V_{bci}$ . The function in Eq. 11 is positive and without poles in the interval  $-\infty < V_{bc} < \infty$ . For the simulators using charge definition, a closed form of the charge can be derived from Eq.11 and combined with Eq.10.

#### Self-heating modeling

Since the temperature coefficients of different model parameters are small (in the order of  $10^{-3}$ ) [8], the changes of the model parameters vs. temperature can be considered linear in a limited temperature range (Eq.12). If required more terms can be added. The values of  $I_c$  and  $V_c$  are available in harmonic balance simulator and are used to calculate dynamically junction temperature  $T_j$ :

$$P_{tot} = I_c V_c + I_b V_b; T_j = P_{tot} R_{therm} + T_{amb} \quad (12)$$

$$I_{jbeT} = I_{jbe}(1 + T_{cljbe}(T_j - T_r)), I_{pkcT} = I_{pkc}(1 + T_{clpc}(T_j - T_r))$$

$$PP_{beT} = PP_{be}(1 + T_{cPPbe}(T_j - T_r))$$

$$V_{jeT} = V_{jeT}(1 + T_c V_{jeT}(T_j - T_r))$$

$$V_{bepT} = V_{bep}(1 + T_c V_{bepT}(T_j - T_r)) \quad \text{where } P_{tot} \text{ is the}$$

dissipated power,  $R_{therm}$  is the thermal resistance. For higher accuracy  $R_{th}$  can be made temperature dependent.

$$R_{thT} = R_{th}(1 + T_{cRth}(T_j - T_r))$$

#### Evaluation of the model

The extraction of parameters starts with the extraction of the DC parameters. When using eq.5.a there are Many fitting programs (Mathematica, MathCAD, Kaleidagraph) can be used to extract basic parameters of the model. Fig.6, 7 show some results from measured and modeled  $I_b$ ,  $I_{ce}$  (measured-points and modeled lines).

There is a large amount of papers on extraction of a small signal equivalent circuit [14-26] of HBT's. The model was implemented as a SDD in MDS (Agilent) with a self-heating and delay part and was experimentally evaluated using DC, S-parameter and Power Spectrum Measurements.

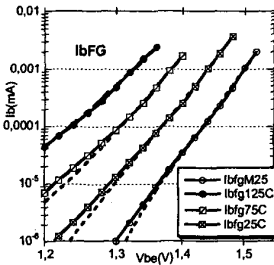


Fig. 6. Meas. and modeled  $I_b$

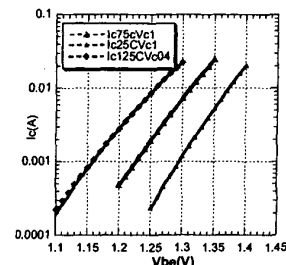


Fig. 7. Meas. and modeled  $I_{ce}$

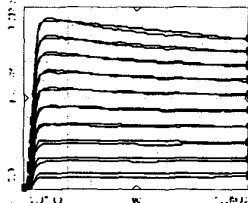


Fig. 8.  $I_{ce}$ :  $I_b = 20 \rightarrow 200 \mu A$  step  $20 \mu A$ . Fig. 9.  $I_{ce}$ ,  $V_{be} 1.28 \rightarrow 1.4$

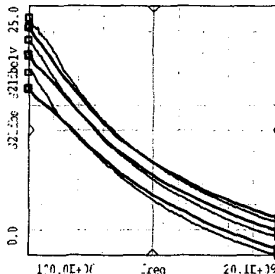
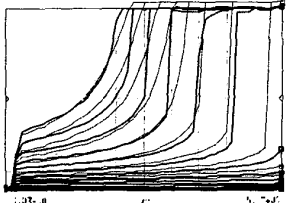


Fig. 10. S21  $V_{be}$ : 1.36- 1.38V

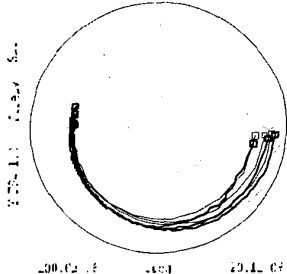


Fig. 11. S11,  $V_{be}$ : 1.36- 1.38V.

In Fig.8, 9 are shown measured and modeled  $I_{ce}$ . In the example shown in Fig. 8 the base current  $I_b$  is a parameter and in Fig. 9  $V_{be}$  is a parameter. The fit is good even with such a simple definition of the model and using only 2 terms in the power series for  $I_{ce}$ . Despite the simplicity the model is able to describe the thermal runaway, Fig. 9.

Fig. 10,11 shows some results of S-parameter measurements and simulations. The model accurately describes the small signal behavior.

The large signal properties of the model were evaluated using a power spectrum (PS) method [28]. Fig. 12 shows results of the measurements and simulations and the fit is good.

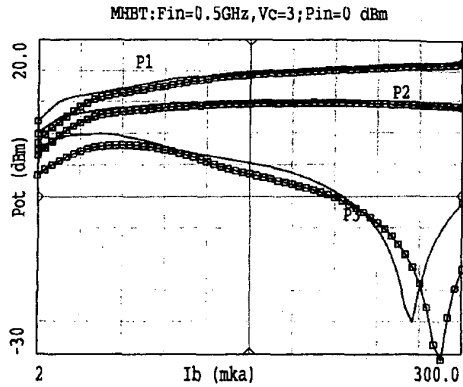


Fig. 12. Measured and modeled PS,  $V_{ce}=3V, I_b \rightarrow 300 \mu A$ .

### III. CONCLUSIONS

A simple bias dependent HBT model applicable to CAD was proposed and implemented. The model was experimentally evaluated with a DC; S-parameter and PS measurements and good correspondence was obtained between the measurements and the model.

### ACKNOWLEDGEMENT

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Table1

$I_{pk}$ mA	$P_{cf1}$	$P_{cf2}$	$V_{bep}$ V	$D_{vpkc}$ V	$\alpha_r$	$\alpha_s$	$B_1$	$B_2$
5.3	28.2	-30	1.31	0.002	18	0.1	0.1	20

$I_{jb}$ mA	$P_{be1}$	$P_{be2}$	$P_{be3}$	$V_{jbe}$	$P_{bc1}$	$V_{jbc}$
1.31	29.8	0	0	1.45	29.8	1.45